

1. Prove These Two Angular Momentum Raising/lowering Operator Relations: $J_+ |j, m\rangle = (j - m)(j + m + 1) |j, m + 1\rangle$

1. Prove These Two Angular Momentum Raising/lowering Operator Relations: $J_- |j, m\rangle = (j + m)(j - m + 1) |j, m - 1\rangle$

Angular momentum in quantum mechanics is a fundamental concept, especially when dealing with systems possessing rotational symmetry. The ladder operators J_+ and J_- play a crucial role in manipulating angular momentum eigenstates, allowing transitions between states with different magnetic quantum numbers m . The relation $J_+ |j, m\rangle = (j - m)(j + m + 1) |j, m + 1\rangle$ encapsulates the action of the raising operator J_+ on the eigenstates $|j, m\rangle$. In this article, we will rigorously derive this relation, explore its significance, and understand the underlying algebraic structure that leads to it.

Understanding the Foundations of Angular Momentum in Quantum Mechanics

The Angular Momentum Operators

In quantum mechanics, angular momentum operators $\hat{J}_x, \hat{J}_y, \hat{J}_z$ satisfy the following commutation relations:

$$[\hat{J}_i, \hat{J}_j] = i\hbar \epsilon_{ijk} \hat{J}_k$$

From these, the total angular momentum operator $\hat{J}^2$ and the z-component $\hat{J}_z$ form a complete set of commuting observables:

$$\begin{aligned} \hat{J}^2 |j, m\rangle &= \hbar^2 j(j+1) |j, m\rangle \\ \hat{J}_z |j, m\rangle &= \hbar m |j, m\rangle \end{aligned}$$

Here, j and m are quantum numbers with $j \geq 0$ and $m \in \{-j, -j+1, \dots, j-1, j\}$.

The Ladder (Raising and Lowering) Operators

Define the ladder operators as:

$$\hat{J}_{\pm} = \hat{J}_x \pm i \hat{J}_y$$

These operators satisfy the following key properties:

- $[\hat{J}_z, \hat{J}_{\pm}] = \pm \hbar \hat{J}_{\pm}$
- $[\hat{J}_+, \hat{J}_-] = 2 \hbar \hat{J}_z$

When acting on the eigenstates $|j, m\rangle$, the ladder operators change the magnetic quantum number m by one unit:

$$\hat{J}_{\pm} |j, m\rangle = c_{\pm} |j, m \pm 1\rangle$$

where $c_{\pm}$ are scalar coefficients to be determined.

Deriving the Action of $\hat{J}_+$ on $|j, m\rangle$

Step 1: Expressing $\hat{J}_+$ Action in Terms of Coefficients

Assuming that $\hat{J}_+$ acts as follows:

$$\hat{J}_+ |j, m\rangle = c_{j,m}^+ |j, m+1\rangle$$

Our goal is to find an explicit expression for $c_{j,m}^+$. To do this, we leverage the properties of the operators and the eigenstates.

Step 2: Use the Hermitian Conjugate Relations

Since $(\hat{J}_+)$ is the Hermitian adjoint of $(\hat{J}_-)$, and the states are orthonormal, the coefficients satisfy certain relations. We can determine $(c_{j,m}^+)$ by considering the norms of the states after the operation.

Step 3: Calculate $(\langle j,m | \hat{J}_- \hat{J}_+ | j,m \rangle)$

Recall that the action of $(\hat{J}_-)$ on $(|j,m+1\rangle)$ is analogous to $(\hat{J}_+)$ but in the opposite direction:

$$\hat{J}_- |j,m+1\rangle = c_{j,m+1}^- |j,m\rangle$$

Using the Hermitian conjugate relation, $(c_{j,m}^+ = c_{j,m+1}^-)$. Therefore, calculating the expectation value:

$$\langle j,m | \hat{J}_- \hat{J}_+ | j,m \rangle = |c_{j,m}^+|^2$$

will help us find the explicit form of $(c_{j,m}^+)$.

Step 4: Express $(\hat{J}_- \hat{J}_+)$ in Terms of $(\hat{J}^2)$ and $(\hat{J}_z)$

Using the angular momentum algebra, note that:

$$\hat{J}_- \hat{J}_+ = \hat{J}^2 - \hat{J}_z^2 - \hbar \hat{J}_z$$

This relation is obtained from the commutation relations and the definitions of ladder operators.

Step 5: Calculate the Expectation Value

Applying this to the state $(|j,m\rangle)$, we get:

$$\langle j, m | \hat{J}_- \hat{J}_+ | j, m \rangle = \langle j, m | (\hat{J}^2 - \hat{J}_z^2 - \hbar \hat{J}_z) | j, m \rangle$$

Using the eigenvalue equations:

$$\begin{aligned} \hat{J}^2 | j, m \rangle &= \hbar^2 j(j+1) | j, m \rangle \\ \hat{J}_z | j, m \rangle &= \hbar m | j, m \rangle \end{aligned}$$

we find:

$$\langle j, m | \hat{J}_- \hat{J}_+ | j, m \rangle = \hbar^2 j(j+1) - \hbar^2 m^2 - \hbar^2 m$$

Step 6: Express $\langle c_{j,m}^+ \rangle$ in Terms of $\langle j \rangle$ and $\langle m \rangle$

Since $\langle c_{j,m}^+ \rangle^2 = \langle j, m | \hat{J}_- \hat{J}_+ | j, m \rangle$, we have:

$$\langle c_{j,m}^+ \rangle^2 = \hbar^2 [j(j+1) - m(m+1)]$$

which simplifies to:

$$\langle c_{j,m}^+ \rangle^2 = \hbar^2 (j - m)(j + m + 1)$$

Step 7: Final Expression for $\langle j_+ | j, m \rangle$

Choosing the phase convention such that $\langle c_{j,m}^+ \rangle$ is positive, the ladder operator's action is:

$$j_+ | j, m \rangle = \hbar \sqrt{(j - m)(j + m + 1)} | j, m+1 \rangle$$

Conclusion: The Explicit Raising Operator Relation

Thus, the relation for the action of the raising operator J_+ on the eigenstates $|j, m\rangle$ is:

$$J_+ |j, m\rangle = \hbar \sqrt{(j - m)(j + m + 1)} |j, m + 1\rangle$$

Similarly, the lowering operator acts as:

$$J_- |j, m\rangle = \hbar \sqrt{(j + m)(j - m + 1)} |j, m - 1\rangle$$

These relations are fundamental in the algebra of angular momentum and are extensively used in quantum mechanics, particularly in the addition of angular momenta, spectral analysis, and the construction of basis states.

Additional Notes and Significance

- The coefficients $\sqrt{(j - m)(j + m + 1)}$

Frequently Asked Questions

What is the significance of the angular momentum raising and lowering operators in quantum mechanics?

They are operators that change the magnetic quantum number m of an eigenstate $|j, m\rangle$ by one unit, allowing transitions between different magnetic sublevels, which is essential for understanding angular momentum in quantum systems.

How do the raising and lowering operators, J_+ and J_- , act on the angular momentum eigenstates $|j, m\rangle$?

They act as ladder operators: $J_+|j, m\rangle = \sqrt{j(j+1) - m(m+1)} |j, m+1\rangle$ and $J_-|j, m\rangle = \sqrt{j(j+1) - m(m-1)} |j, m-1\rangle$, raising or lowering the m value by one.

What is the derivation of the relation $J_+|j, m\rangle = (j-m) |j, m+1\rangle$?

This relation follows from the algebra of angular momentum operators and the normalization of eigenstates, specifically using the commutation relations and the action of J_+ on $|j, m\rangle$ to find the proportionality constant, which is $(j-m)$.

Can you explain the mathematical steps to prove that $J_+|j, m\rangle = (j-m) |j, m+1\rangle$?

Yes. Starting from the eigenvalue equations and the commutation relations, you apply J_+ to $|j, m\rangle$, then compute the inner product with $|j, m+1\rangle$, and use normalization to find the coefficient as $\sqrt{j(j+1) - m(m+1)}$. Recognizing that this equals $(j-m)$, the relation is established.

Why is it important to verify the relation $J_+|j, m\rangle = (j-m) |j, m+1\rangle$?

Verifying this relation confirms the proper action of ladder operators on angular momentum eigenstates, which is fundamental for calculations involving transitions, selection rules, and spectral lines in quantum mechanics.

Are these operator relations specific to certain values of j and m ?

No, these relations are general for all valid quantum numbers j and m , where j is the total angular momentum quantum number and m ranges from $-j$ to j in integer steps.

What role do these raising/lowering operator relations play in angular momentum coupling?

They facilitate the addition of angular momenta by allowing the systematic construction of combined eigenstates and the calculation of Clebsch-Gordan coefficients, which describe how individual angular momenta combine.

How can these relations be applied in practical quantum mechanics problems?

They are used to determine transition probabilities, compute matrix elements, analyze spectral lines, and solve problems involving atomic, molecular, and nuclear angular momentum states.

Additional Resources

Prove These Two Angular Momentum Raising/lowering Operator Relations: $\hat{J}_+ |j, m\rangle = (j - m) |j, m+1\rangle$ and $\hat{J}_- |j, m\rangle = (j + m) |j, m-1\rangle$

Introduction

In quantum mechanics, the angular momentum operators are fundamental to understanding the rotational symmetries and the quantization of angular momentum in atomic, molecular, and subatomic systems. Central to this framework are the ladder (or raising and lowering) operators, $\hat{J}_+$ and $\hat{J}_-$, which serve to elevate or depress the magnetic quantum number (m) associated with a total angular momentum quantum number (j) .

The relations:

$$\begin{aligned} \hat{J}_+ |j, m\rangle &= \sqrt{(j - m)(j + m + 1)} |j, m+1\rangle, \\ \hat{J}_- |j, m\rangle &= \sqrt{(j + m)(j - m + 1)} |j, m-1\rangle, \end{aligned}$$

are foundational in the algebraic treatment of angular momentum. They not only facilitate the calculation of matrix elements but also underpin the structure of the representation theory of the Lie algebra $\mathfrak{su}(2)$.

In this comprehensive review, we aim to rigorously derive and prove the simplified relations:

$$\begin{aligned} j_{+} |j, m\rangle &= (j - m) |j, m+1\rangle, \\ j_{-} |j, m\rangle &= (j + m) |j, m-1\rangle, \end{aligned}$$

which are often presented in textbooks under specific normalization conventions.

Background and Mathematical Framework

The Angular Momentum Algebra

The angular momentum operators $(\hat{J}_x, \hat{J}_y, \hat{J}_z)$ satisfy the commutation relations:

$$[\hat{J}_i, \hat{J}_j] = i \hbar \epsilon_{ijk} \hat{J}_k,$$

where (ϵ_{ijk}) is the Levi-Civita symbol.

Defining the ladder operators:

$$\hat{J}_{\pm} = \hat{J}_x \pm i \hat{J}_y,$$

these satisfy:

$$\begin{aligned} [\hat{J}_z, \hat{J}_{\pm}] &= \pm \hbar \hat{J}_{\pm}, \\ [\hat{J}_+, \hat{J}_-] &= 2 \hbar \hat{J}_z. \end{aligned}$$

The simultaneous eigenstates $(|j, m\rangle)$ are characterized by:

$$\begin{aligned} \hat{J}^2 |j, m\rangle &= \hbar^2 j(j+1) |j, m\rangle, \\ \hat{J}_z |j, m\rangle &= \hbar m |j, m\rangle, \end{aligned}$$

where (j) is a non-negative half-integer or integer, and (m) takes on values $(-j, -j+1, \dots, j-1, j)$.

Normalization Convention

Standard conventions often normalize the states such that:

$$\langle j, m | j, m \rangle = 1,$$

and the ladder operators act as:

$$\hat{J}_{\pm} |j, m\rangle = \hbar \sqrt{(j \mp m)(j \pm m + 1)} |j, m \pm 1\rangle.$$

However, in certain derivations or textbooks, the states are scaled such that the relation simplifies, leading to the relations under investigation.

Derivation of the Ladder Operator Relations

Step 1: Action of $\hat{J}_+$ and $\hat{J}_-$ on $|j, m\rangle$

The key starting point is the action of the ladder operators:

$$\hat{J}_{\pm} |j, m\rangle = c_{\pm}(j, m) |j, m \pm 1\rangle,$$

where $c_{\pm}(j, m)$ are coefficients to be determined.

The goal is to find explicit forms of $c_{\pm}(j, m)$ that satisfy the algebra and normalization conditions.

Step 2: Use of Commutation Relations and Casimir Operator

The Casimir operator:

$$\hat{J}^2 = \hat{J}_x^2 + \hat{J}_y^2 + \hat{J}_z^2,$$

commutes with all components and acts on the states as:

$$\hat{J}^2 |j, m\rangle = \hbar^2 j(j+1) |j, m\rangle.$$

Applying $\hat{J}_{\pm}$ multiple times and using the commutation relations yields recurrence relations for the coefficients.

Detailed Proof of the Relations

Step 3: Establishing the Coefficients

Starting from the known algebraic relations, the coefficients are deduced:

$$c_{\pm}(j,m) = \hbar \sqrt{(j \mp m)(j \pm m + 1)}.$$

This derivation hinges on the normalization of states and the algebraic properties of the operators.

Explicit derivation:

1. Consider the norm of the state $|\hat{J}_{\pm}|j,m\rangle$:

$$\langle j,m | \hat{J}_{\pm} \hat{J}_{\pm} | j,m \rangle = |c_{\pm}(j,m)|^2,$$

since the states are normalized.

2. Use the commutation relation:

$$[\hat{J}_{\pm}, \hat{J}_{\pm}] = \hat{J}^2 - \hat{J}_z^2 \mp \hbar \hat{J}_z,$$

which follows from the algebra.

3. Acting on $|j,m\rangle$:

$$\langle j,m | \hat{J}_{\pm} \hat{J}_{\pm} | j,m \rangle = [\hbar^2 j(j+1) - \hbar^2 m^2 \mp \hbar m].$$

4. Simplify to obtain:

$$|c_{\pm}(j,m)|^2 = \hbar^2 [j(j+1) - m(m \pm 1)].$$

5. Taking the square root:

$$c_{\pm}(j,m) = \hbar \sqrt{(j \mp m)(j \pm m + 1)}.$$

Simplification and Normalization Choices

Standard Convention

Choosing units such that $\hbar = 1$, the relations become:

$$\begin{aligned} j_{+} |j, m\rangle &= \sqrt{(j-m)(j+m+1)} |j, m+1\rangle, \\ j_{-} |j, m\rangle &= \sqrt{(j+m)(j-m+1)} |j, m-1\rangle. \end{aligned}$$

Simplified Relations (Given in the Problem)

The relations to be proved are:

$$\begin{aligned} j_{+} |j, m\rangle &= (j-m) |j, m+1\rangle, \\ j_{-} |j, m\rangle &= (j+m) |j, m-1\rangle, \end{aligned}$$

which correspond to a particular normalization of the states where the coefficients are linear functions of j and m , rather than square roots.

Proof of the Simplified Relations

Assumption: Normalization of States

Suppose the states are normalized such that:

$$\begin{aligned} j_{+} |j, m\rangle &= (j-m) |j, m+1\rangle, \\ j_{-} |j, m\rangle &= (j+m) |j, m-1\rangle, \end{aligned}$$

and the states $|j, m\rangle$ form an orthonormal basis.

Step 1: Show that the relations are compatible with the algebra

1. From the given relations, compute the commutator:

$j_{+} j_{-}$

$$[j_+, j_-] |j, m\rangle = (j_+ j_- - j_- j_+) |j, m\rangle.$$

2. Applying the relations:

$$j_+ j_- |j, m\rangle = j_+ (j + m) |j, m-1\rangle = (j + m) j_+ |j, m-1\rangle,$$

$$j_+ |j, m-1\rangle = (j - (m-1)) |j, m\rangle = (j - m + 1) |j, m\rangle,$$

thus:

$$j_+ j_- |j, m\rangle = (j + m) (j - m + 1) |j, m\rangle.$$

Similarly:

$$j_- j_+ |j, m\rangle = (j - m) (j + m + 1) |j, m\rangle.$$

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